

Six-dimensional $N = 2$, $F_4/SO(9)$ Supergravity and the roof of the Magic Square

L. J. Romans

Larry.J.Romans@gmail.com

Abstract

Motivated by the construction by Günaydin, Sierra and Townsend of Maxwell-Einstein $N = 2$ supergravity theories in 3, 4 and 5 dimensions realizing symmetries appearing in the lower three rows of the 4×4 Freudenthal-Tits Magic Square, we present a six-dimensional theory with $F_4/SO(9)$ structure, corresponding to the maximal ‘‘octonionic’’ item in the top row of the square. The construction involves detailed interplay between the scalar coset structure and the spin-1 sector comprised of vectors, self-dual and anti-self-dual antisymmetric tensors.

1. Introduction

The celebrated Magic Square of Freudenthal and Tits

	$\mathbb{A} = \mathbb{O}$	$\mathbb{A} = \mathbb{H}$	$\mathbb{A} = \mathbb{C}$	$\mathbb{A} = \mathbb{R}$
$d = 6$	F_4	C_3	A_2	A_1
$d = 5$	E_6	A_5	$A_2 \times A_2$	A_2
$d = 4$	E_7	D_6	A_5	C_3
$d = 3$	E_8	E_7	E_6	F_4

	$\mathbb{A} = \mathbb{O}$	$\mathbb{A} = \mathbb{H}$	$\mathbb{A} = \mathbb{C}$	$\mathbb{A} = \mathbb{R}$
$d = 6$	$\frac{F_{4(-20)}}{SO(9)}$	$\frac{USp(4, 2)}{USp(4) \times SU(2)}$	$\frac{SU(2, 1)}{SU(2) \times U(1)}$	$\frac{SO(2, 1)}{SO(2)}$
$d = 5$	$\frac{E_{6(-26)}}{F_4}$	$\frac{SU^*(6)}{USp(6)}$	$\frac{SL(3, \mathbb{C})}{SU(3)}$	$\frac{SL(3, \mathbb{R})}{SO(3)}$
$d = 4$	$\frac{E_{7(-25)}}{E_6 \times U(1)}$	$\frac{SO^*(12)}{SU(6) \times U(1)}$	$\frac{SU(3, 3)}{SU(3) \times SU(3) \times U(1)}$	$\frac{USp(6)}{SU(3) \times U(1)}$
$d = 3$	$\frac{E_{8(-24)}}{E_7 \times SU(2)}$	$\frac{E_{7(-5)}}{SO(12) \times SU(2)}$	$\frac{E_{6(2)}}{SU(6) \times SU(2)}$	$\frac{F_{4(4)}}{USp(6) \times SU(2)}$

$d = 6$	$\frac{26 = \underline{1}^{(t+)} + \underline{16}^{(t-)} + \underline{9}^{(v)}}{52 - 36 = 16^{(0)}}$	$\frac{14 = \underline{1}^{(t+)} + \underline{8}^{(t-)} + \underline{5}^{(v)}}{21 - (10 + 3) = 8^{(0)}}$	$\frac{8 = \underline{1}^{(t+)} + \underline{4}^{(t-)} + \underline{3}^{(v)}}{8 - (3 + 1) = 4^{(0)}}$	$\frac{5 = \underline{1}^{(t+)} + \underline{2}^{(t-)} + \underline{2}^{(v)}}{3 - 1 = 2^{(0)}}$
$d = 5$	$\frac{27 = \underline{1}^{(v)} + \underline{26}^{(v)}}{78 - 52 = 26^{(0)}}$	$\frac{15 = \underline{1}^{(v)} + \underline{14}^{(v)}}{35 - 21 = 14^{(0)}}$	$\frac{9 = \underline{1}^{(v)} + \underline{8}^{(v)}}{16 - 8 = 8^{(0)}}$	$\frac{6 = \underline{1}^{(v)} + \underline{5}^{(v)}}{8 - 3 = 5^{(0)}}$
$d = 4$	$\frac{28 = \underline{1}^{(v)} + \underline{27}^{(v)}}{133 - (78 + 1) = 54^{(0)}}$	$\frac{16 = \underline{1}^{(v)} + \underline{15}^{(v)}}{66 - (35 + 1) = 30^{(0)}}$	$\frac{10 = \underline{1}^{(v)} + \underline{9}^{(v)}}{35 - (8 + 8 + 1) = 18^{(0)}}$	$\frac{7 = \underline{1}^{(v)} + \underline{6}^{(v)}}{21 - (8 + 1) = 12^{(0)}}$
$d = 3$	$248 - (133 + 3) = 112^{(0)}$	$133 - (66 + 3) = 64^{(0)}$	$78 - (35 + 3) = 40^{(0)}$	$52 - (21 + 3) = 28^{(0)}$

2. $F_{4(-20)}$ description in an $SO(9)$ basis

$$\begin{pmatrix} Z_0 & \underline{1} \\ Z_i & \underline{9} \\ Z_a & \underline{16} \end{pmatrix}$$

$$\begin{aligned} \mathcal{Q} &\equiv \eta^{IJ} Z_I Z_J \\ &= Z_0^2 + Z_i Z_i \mp Z_a Z_a \end{aligned} \quad (1)$$

$$\begin{aligned} \mathcal{C} &\equiv c^{IJK} Z_I Z_J Z_K \\ &= Z_0^3 - 3Z_0 Z_i Z_i \mp \frac{3}{2} Z_0 Z_a Z_a \pm \frac{3\sqrt{3}}{2} (\Gamma_i^{ab}) Z_i Z_a Z_b \end{aligned} \quad (2)$$

$$\begin{pmatrix} 0 & 0 & \sqrt{3}Y^b \\ 0 & X^{ij} & (\Gamma_i^{bc}) Y^c \\ \sqrt{3}Y^a & (\Gamma_j^{ac}) Y^c & \frac{1}{4} (\Gamma_{kl}^{ab}) X^{kl} \end{pmatrix} \quad (3)$$

2. The scalar sector

$$(v^I_0 \quad v^I_i \quad v^I_a) \begin{pmatrix} u^0_J \\ u^i_J \\ u^a_J \end{pmatrix} = (\delta^I_J) \quad (4)$$

$$\begin{pmatrix} u^0_I \\ u^i_I \\ u^a_I \end{pmatrix} (v^I_0 \quad v^I_j \quad v^I_b) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \delta^i_j & 0 \\ 0 & 0 & \delta^a_b \end{pmatrix} \quad (5)$$

$$\begin{pmatrix} u^0_I \\ u^i_I \\ u^a_I \end{pmatrix} (\partial_\mu v^I_0 \quad \partial_\mu v^I_j \quad \partial_\mu v^I_b) = \begin{pmatrix} 0 & 0 & \sqrt{3}R_\mu^b \\ 0 & Q_\mu^{ij} & (\Gamma_i^{bc}) R_\mu^c \\ \sqrt{3}R_\mu^a & (\Gamma_j^{ac}) R_\mu^c & \frac{1}{4} (\Gamma_{kl}^{ab}) Q_\mu^{kl} \end{pmatrix} \quad (6)$$

$$\begin{pmatrix} u^0_I \\ u^i_I \\ u^a_I \end{pmatrix} (\mathcal{D}_\mu v^I_0 \quad \mathcal{D}_\mu v^I_j \quad \mathcal{D}_\mu v^I_b) = \begin{pmatrix} 0 & 0 & \sqrt{3}R_\mu^b \\ 0 & 0 & (\Gamma_i^{bc}) R_\mu^c \\ \sqrt{3}R_\mu^a & (\Gamma_j^{ac}) R_\mu^c & 0 \end{pmatrix} \quad (7)$$

$$\mathcal{D}_\mu v^I_0 \equiv \partial_\mu v^I_0 \quad (8)$$

$$\mathcal{D}_\mu v^I_i \equiv \partial_\mu v^I_i - Q_\mu^{ij} v^I_j \quad (9)$$

$$\mathcal{D}_\mu v^I_a \equiv \partial_\mu v^I_a - \frac{1}{4} Q_\mu^{ij} (\Gamma_{ij}^{ab}) v^I_b \quad (10)$$

The $SO(9)$ Fierz identities

$$\Gamma_i^{a(b} \Gamma_i^{cd)} = \delta^{a(b} \delta^{cd)} \quad (11)$$

$$3\delta^{a[c} \delta^{d]b} + \Gamma_i^{a[c} \Gamma_i^{d]b} = \frac{1}{4} \Gamma_{ij}^{ab} \Gamma_{ij}^{cd} \quad (12)$$

3. The spin-1 sector

$$F_{\mu\nu}^I = 2\partial_{[\mu}A_{\nu]}^I \quad (13)$$

$$G_{\mu\nu\rho}^I = 3\partial_{[\mu}B_{\nu\rho]}^I + 3c_{JK}^I F_{[\mu\nu}^J A_{\rho]}^K \quad (14)$$

$$\delta A_\mu^I = \partial_\mu \lambda^I \quad (15)$$

$$\delta B_{\mu\nu}^I = 2\partial_{[\mu}\Lambda_{\nu]}^I + 3c_{JK}^I F_{\mu\nu}^J \lambda^K \quad (16)$$

$$\partial_{[\mu}F_{\nu\rho]}^I = 0 \quad (17)$$

$$\partial_{[\mu}G_{\nu\rho\sigma]}^I = 3c_{JK}^I F_{[\mu\nu}^J F_{\rho\sigma]}^K \quad (18)$$

$$F_{\mu\nu}^0 \equiv u^0{}_I F_{\mu\nu}^I \quad (19)$$

$$F_{\mu\nu}^i \equiv u^i{}_I F_{\mu\nu}^I \quad (20)$$

$$F_{\mu\nu}^a \equiv u^a{}_I F_{\mu\nu}^I \quad (21)$$

$$F_{\mu\nu}^0 = 0 \quad (22)$$

$$F_{\mu\nu}^a = 0 \quad (23)$$

$$G_{\mu\nu\rho}^0 \equiv u^0{}_I G_{\mu\nu\rho}^I \quad (24)$$

$$G_{\mu\nu\rho}^i \equiv u^i{}_I G_{\mu\nu\rho}^I \quad (25)$$

$$G_{\mu\nu\rho}^a \equiv u^a{}_I G_{\mu\nu\rho}^I \quad (26)$$

$$G_{\mu\nu\rho}^0 = +\tilde{G}_{\mu\nu\rho}^0 \quad (27)$$

$$G_{\mu\nu\rho}^a = -\tilde{G}_{\mu\nu\rho}^a \quad (28)$$

$$G_{\mu\nu\rho}^i = 0 \quad (29)$$

4. Supersymmetry transformations and field equations

$$\delta e_{\mu\alpha} = \bar{\psi}_\mu \gamma^\alpha \epsilon \quad (30)$$

$$\delta \psi_\mu = \nabla_\mu \epsilon + \frac{1}{4} G_{\mu\nu\rho}^0 \gamma^{\nu\rho} \epsilon \quad (31)$$

$$\delta A_\mu^I = v^I{}_i \bar{\lambda}^i \gamma_\mu \epsilon \quad (32)$$

$$\delta B_{\mu\nu}^I = v^I{}_0 \bar{\psi}_{[\mu} \gamma_{\nu]} \epsilon - \frac{1}{2} v^I{}_a \bar{\chi}^a \gamma_{\mu\nu} \epsilon + 2c_{JK}^I A_{[\mu}^J \delta A_{\nu]}^K \quad (33)$$

$$\delta \lambda^i = -\frac{1}{4} F_{\mu\nu}^i \gamma^{\mu\nu} \epsilon \quad (34)$$

$$\delta \chi^a = P_\mu^a \gamma^\mu \epsilon + \frac{1}{12} G_{\mu\nu\rho}^a \gamma^{\mu\nu\rho} \epsilon \quad (35)$$

$$[\delta_{s.s.}(\epsilon_2), \delta_{s.s.}(\epsilon_1)] = \delta_{g.c.}(\xi^\mu) + \delta_{l.l.}(\Sigma^{\alpha\beta}) \quad (36)$$

$$+ \delta_{gauge(v)}(\lambda^I) + \delta_{gauge(t)}(\Lambda_\mu^I) \quad (37)$$

$$+ \delta_{SO(9)}(L^{ij}) \quad (38)$$

$$\xi^\mu = \bar{\epsilon}_2 \gamma^\mu \epsilon_1 \quad (39)$$

$$\Sigma^\mu = \xi^\mu (\omega_\mu^{\alpha\beta} - G^0{}_\mu{}^{\alpha\beta}) \quad (40)$$

$$\lambda^I = -\xi^\mu A_\mu^I \quad (41)$$

$$\Lambda_\mu^I = \dots \quad (42)$$

$$L^{ij} = -\xi^\mu Q_\mu^{ij} \quad (43)$$

5. Discussion

A. Spacetime conventions

$$\tilde{X}_{\mu\nu\rho} = \frac{1}{6} e_{\mu\nu\rho}{}^{\sigma\tau\kappa} X_{\sigma\tau\kappa} \quad (44)$$

References

- [1] M. Günaydin, G. Sierra and P. K. Townsend, Nucl. Phys. B242 (1984) 244.
- [2] H. Freudenthal, Nederl. Akad. Wetensch. Proc. Ser. A, 57 (1954) 218; T. A. Springer, ibid. 65 (1962) 259; J. Tits, ibid. 65 (1962) 530; 69 (1966) 223.
- [3] L. J. Romans, Nucl. Phys. B276 (1986) 71.
- [4] J. H. Schwarz, Nucl. Phys. B226 (1983) 269.